# **Recent developments of molecular dynamics calculation software MODYLAS**

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# Introduction



NU has six Nobel Prize winners:

Chemistry R. Noyori [2001], O. Shimomura [2008] Physics M. Kobayashi & T. Maskawa [2008], I. Akasaki & H. Amano [2014]

# Introduction

#### Department of Materials Chemistry, Nagoya University

• <u>Theoretical and Computational Chemistry</u> Other groups:

- Catalyst Design
- Material Design Chemistry
- Energy Conversion Chemistry

## Staff of Okazaki Lab.

Professor: Susumu Okazaki

Associate Professor: Wataru Shinoda Assistant Professor: Kazushi Fujimoto



Designated Associate Prof.: Noriyuki Yoshii (CCS)  $\rightarrow$  Theory of FMM Designated Associate Prof.: Yoshimichi Andoh (CCS)  $\rightarrow$  Implementation & parallelization Dr. Tatsuya Sakashita  $\rightarrow$  Implementation of solid harmonics, etc.

, and 8 post-docs; 3 Ph.D, 12 Master course, 5 undergraduate students

#### **Research Subjects**

- Plasma membranes of normal and cancerous cells
- Viruses, and their recognition by receptors
- Micelle formation and micellar solubilization
- Polymer material
- Development of theory of MD calculation, and MD simulation software, MODYLAS



粗視化シミュレーションによるマルチラメラベシクルの生成





マルチラメラベシクル、小児麻痺ウイルスの分子動力学シミュレーション

# **Molecular dynamics (MD) calculation** 4/27

**Basic equation** : Newton's equation of motion

$$m_{i} \frac{d^{2} \mathbf{r}_{i}}{dt^{2}} = \mathbf{F}_{i}$$

$$i = 1, \dots, N$$
Numerical
integration

*t* : time

 $\mathbf{r}_i$ : atom coordinate

 $F_i$ : forces working on each atom

 $m_i$ : atom mass

 $\dot{N}$ : number of atoms O(10<sup>4</sup>-10<sup>7</sup>)

2nd order ordinary differential equations for 3N degree of freedoms [initial conditions:  $r_i(0), v_i(0)$ ] Generalization Verlet's method  $\longrightarrow$  RESPA[1]  $v_i\left(t + \frac{\Delta t}{2}\right) = v_i(t) + \frac{\Delta t}{2} \frac{F_i(\Delta t)}{m_i}$   $r_i(t + \Delta t) = r_i(t) + \Delta t v_i\left(t + \frac{\Delta t}{2}\right)$  cycle  $v_i\left(t + \Delta t\right) = v_i\left(t + \frac{\Delta t}{2}\right) + \frac{\Delta t}{2} \frac{F_i(t + \Delta t)}{m_i}$ 

 $\boldsymbol{v}_i$ : atom velocitiy

 $\Delta t$ : time step size O(10<sup>-15</sup>) sec

100 ns long MD calculations =  $10^8$  cycles

[1] M. Tuckermann et al. J. Chem. Phys. 97, 1990 (1992).

## **Potential function**

$$\boldsymbol{F}_{i} = -\frac{\partial \boldsymbol{\Phi}_{tot}(\boldsymbol{r}_{1}, \dots, \boldsymbol{r}_{N})}{\partial \boldsymbol{r}_{i}} \qquad \boldsymbol{\Phi}_{tot}(\boldsymbol{r}_{1}, \dots, \boldsymbol{r}_{N}) : \text{Total potential energy}$$
(Classical model)



## **Review of FMM**

• O(N) algorithm

[2] L.F. Greengard (1987).

• suit for massive parallelization

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## **Accuracy of our coded FMM**

The deviation of a force from the correct value obtained by the Ewald method:

$$\sigma_{\alpha} = \sqrt{\frac{\sum_{i}^{N} \left(F_{\alpha,i}^{P} - F_{\alpha,i}^{\text{Ewald}}\right)^{2}}{\sum_{i}^{N} \left(F_{\alpha,i}^{\text{Ewald}}\right)^{2}}} \quad \alpha = x, y, z$$

*P*: The order of expansion, noted as "nmax" in the following slides

$$\varPhi_{\text{far}} = \sum_{n=0}^{P} \sum_{m=-n}^{n} L_n^m Y_n^m(\theta, \phi) r^n$$

Deviation  $\sigma$  (=  $\sigma_x = \sigma_y = \sigma_z$ ) of the forces calculated by the FMM from the correct ones by the Ewald method.

[3] Y. Andoh, N. Yoshii, K. Fujimoto, et al., "MODYLAS: A highly parallelized general-purpose molecular dynamics simulation program for large-scale systems with long-range forces calculated by fast multipole method (FMM) and highly scalable fine-grained new paral processing algorithms", J. Chem. Theory Comput., **2013**, 9, 3201.



## Strong scaling test on the K computer

Measurement in March, 2012



# **Fig.** Measured overall calculation time per MD step ( $\Delta$ t) with respect to the 64-node calculation.

# Fig. Measured partial calculation times and communication time per MD step $(\Delta t)$

[3] Y. Andoh, N. Yoshii, K. Fujimoto, et al., "MODYLAS: A highly parallelized general-purpose molecular dynamics simulation program for large-scale systems with long-range forces calculated by fast multipole method (FMM) and highly scalable fine-grained new paral processing algorithms", J. Chem. Theory Comput., **2013**, 9, 3201.

# **Applications of MODYLAS**





All-atomistic MD simulation of entire poliovirus empty capsid in solution *J. Chem. Phys.* **141**, 165101 (2014).

Examples of collaboration with companies

- Nippon-Kayaku J. Phys. Chem. B, 113, 15181 (2009). Micelle of PEG-PBLG-Ac block copolymers
  - 0.2 M atoms 7ns

• Dainihon Sumitomo Seiyaku PLoS ONE, **11**, 1 (2016).

Drug-membrane protein affinity

Nitto Denko
 Polymer adhesives



- S.Okazaki, H.Tanaka,
- *J. Phys. Chem. B*, **118**, 1900 (2014). *J. Phys. Chem. B*, **118**, 11797 (2014).;

















ImPACT Project (Prof. Ito): Realizing Ultra-Thin and Flexible Tough Polymers

MODYLAS is not benchmark software, but practical general-purpose MD software.

革新的研究開発推進プログラム

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- Molecular dynamics calculation combined with FMM
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## **Functions:**

FMM for rectangular unit cell, with anisotropic cell division Pressure tensor calculated by the FMM (J. Comput. Chem., in press) Implementation of the regularized FMM Implementation of bond-breaking potential for polymer simulation Implementation of the EBV potential for proton-transfer simulation

Conclusion

Red: today's topics

## Performance: 1. Improvement thread-level parallelization (p2p)



## 2. Replacement of the spherical harmonics by solid harmonies<sup>27</sup>

P2M (Particle to Multipole, calculation of multipole moments)

$$M_{\ell}^{m} = \sum_{i=1}^{N} q_{i} R_{\ell}^{m} (\mathbf{r}_{M} - \mathbf{r}_{i})$$
Regular solid harmonics
$$R_{\ell}^{m}(r, \theta, \phi) = \frac{r^{\ell}}{(\ell + m)!} P_{\ell}^{m}(\cos \theta) e^{im\varphi}$$

M2M (Multipole to Multipole, to translate of expansion center of M)

$$M'_{\ell}^{m} = \sum_{\lambda=0}^{n_{\max}} \sum_{\mu=-\lambda}^{\lambda} M_{\ell-\lambda}^{m-\mu} R_{\lambda}^{\mu} (\mathbf{r}_{M'} - \mathbf{r}_{M})$$
Traditional M2L with  $Y_{n}^{m}$ 

$$L_{n}^{m} = \sum_{j=0}^{P} \sum_{k=-j}^{j} \frac{J_{m}^{k} M_{j}^{k} A_{j}^{k} A_{n}^{m} Y_{n+j}^{k-m} (\alpha, \mu)}{A_{n+j}^{k-m} \rho^{n+j+1}}$$
M2L (Multipole to Local expansion coeffs)

$$L_{\ell}^{m} = \sum_{\lambda=0}^{n_{\max}} \sum_{\mu=-\lambda}^{\lambda} M_{\lambda}^{\mu} S_{\ell+\lambda}^{-(m+\mu)} (\mathbf{r}_{L} - \mathbf{r}_{M})$$
Singular solid harmonics  
$$S_{\ell}^{m}(r, \theta, \phi) = (-1)^{\ell+m} \frac{(\ell-m)!}{r^{\ell+1}} P_{\ell}^{m}(\cos \theta) e^{im\varphi}$$

L2L (Local expansion to Local expansion, to translate of expansion center of L)

$$L'_{\ell}^{m} = \sum_{\lambda=\ell}^{n_{\max}} \sum_{\mu=-\lambda}^{\lambda} L_{\ell}^{m} R_{\lambda-\ell}^{\mu-m} (\mathbf{r}_{L'} - \mathbf{r}_{L})$$
  

$$\Phi_{i}(\mathbf{r}) = q_{i} \sum_{\ell=0}^{n_{\max}} \sum_{m=-\ell}^{\ell} L_{\ell}^{m} R_{\ell}^{m} (\mathbf{r} - \mathbf{r}_{L}) \qquad \mathbf{F}_{i} = -q_{i} \sum_{\ell=0}^{n_{\max}} \sum_{m=-\ell}^{\ell} L_{\ell}^{m} \begin{bmatrix} \frac{1}{2} (R_{\ell-1}^{m+1} - R_{\ell-1}^{m-1}) \\ -\frac{1}{2} (R_{\ell-1}^{m+1} + R_{\ell-1}^{m-1}) \\ R_{\ell-1}^{m} \end{bmatrix}$$

#### Code has become simple greatly, make extension of FMM algorithm much easier

We have succeeded in reducing arithmetic amount of FMM by implementing the following three methods

(1) Utilizing symmetric nature of matrix element of multipole expansion coefficients M and local expansion coefficients  $L^{[4]}$ 

(2) Usage of recurrence formula<sup>[4]</sup> which can avoid a conversion from Cartesian coordinate to polar coordinate.

(3) Rotation of the axes in M2L operation<sup>[5]</sup>

[4] Basic idea is described in unpublished work by K. Nitadori, "Particle mesh multipole method: An efficient solver for gravitational/electrostatic forces based on multipole method and fast convolution over a uniform mesh".

[5] C. A. White, M. Head-Gordon. J. Chem. Phys., Vol. 105, No. 12, pp. 5061–5067, 1996.

### **Extension of function:**

## 1. FMM for rectangular unit cell with anisotropic cell division



Work by Y. Andoh & Prof. N. Yoshii

#### A: Traditional FMM:

- Cubic unit cell
- Isotropic cell division, following binary branches (octree structure)

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• NPROCS =  $2^n$ 

#### B: Our extended FMM<sup>[6]</sup>:

- Rectangular unit cell
- Anisotropic cell division, following binary and ternary branches, independently selectable for each axis and each level
- NPROCS =  $2^{n*}3^m$  is available

[6]Y. Andoh, N. Yoshii, S. Okazaki, *will be submitted soon.* 

## **Extension of function:**

## 2. New expression for pressure tensor in FMM



transform	formula			
P2M	$\frac{\partial M_{s}^{m}(I_{\max})}{\partial h_{\alpha\beta}} = \sum_{i=1}^{N} q_{i} \left[ \nabla_{i} Y_{s}^{-w} \left( \hat{\boldsymbol{\rho}}_{i} \right) \boldsymbol{\rho}_{i}^{n} \right]_{\alpha} s_{i\beta}$			
M2M	$\frac{\partial M_{j}^{k}(l+1)}{\partial h_{\alpha\beta}} = \sum_{n=0}^{j} \sum_{m=n}^{n} T_{\text{MM},j,u}^{k,m} \left\{ \frac{\partial M_{j-n}^{k-m}(l)}{\partial h_{\alpha\beta}} Y_{n}^{-m}(\hat{\boldsymbol{\rho}}_{\text{MM}}) \rho_{\text{MM}}^{m} + M_{j-n}^{k-m}(l) \left[ \nabla_{\text{MM}} Y_{n}^{-m}(\hat{\boldsymbol{\rho}}_{\text{MM}}) \rho_{\text{MM}}^{m} \right]_{\alpha} S_{\text{MM}\beta} \right\}$			
M2L	$\frac{\partial L_{j}^{k}(l)}{\partial h_{\alpha\beta}} = \sum_{n=0}^{\infty} \sum_{m=-n}^{n} T_{\mathrm{M}_{-j,n}}^{k,m} \left\{ \frac{\partial \mathcal{M}_{n}^{m}(l)}{\partial h_{\alpha\beta}} \frac{Y_{n+j}^{m+k}\left(\hat{\rho}_{\mathrm{ML}}\right)}{\rho_{\mathrm{M}_{-}}^{n+j+1}} + \mathcal{M}_{n}^{m}(l) \left[ \nabla_{\mathrm{M}_{-}} \frac{Y_{n+j}^{m-k}\left(\hat{\rho}_{\mathrm{M}_{-}}\right)}{\rho_{\mathrm{M}_{-}}^{n+j+1}} \right]_{\alpha} s_{\mathrm{ML},\beta} \right\}$			
M2L	$\sum_{n=0}^{V_{mm}} \frac{Y_n^m(\hat{\rho}_v)}{\rho_n^{n+1}} = \sum_{n=0}^{V_{mm}} \frac{Y_n^m(\hat{\rho}_v)}{\rho_n^{n+1}} \frac{\Gamma(n+1/2,\kappa^2\rho_v^2)}{\Gamma(n+1/2)} + \sum_{n=0}^{H_{mm}} \frac{i^n \pi^{n-l/2} Y_n^m(\hat{\nu}_H) v_H^{n-2} \exp(-\pi^2 v_H^2 / \kappa^2)}{V\Gamma(n+1/2)}$			
(image cells)				
M2L	$\left \frac{\partial}{\partial h_{off}}\sum_{x=0}^{v_{m}}\frac{Y_{x}^{m}\left(\hat{\rho}_{x}\right)}{\rho_{x}^{n+1}}\frac{\Gamma\left(n+1/2,\kappa^{2}\rho_{x}^{2}\right)}{\Gamma\left(n+1/2\right)}=\sum_{x=0}^{v_{m}}\frac{1}{\Gamma\left(n+1/2\right)}\left[\Gamma\left(n+1/2,\kappa^{2}\rho_{x}^{2}\right)\nabla_{v}\frac{Y_{m}^{m}\left(\hat{\rho}_{v}\right)}{\rho_{v}^{n+1}}-2\frac{\kappa^{n+1}e^{-\kappa^{2}\rho_{x}^{2}}Y_{x}^{m}\left(\hat{\rho}_{v}\right)}{\rho_{v}^{2}}\rho_{v}\right]s_{vff}\right $			
(image cells)	$\frac{\partial h_{\alpha\beta}}{\partial r} \frac{\rho_{\nu}^{n+1}}{\partial r} = \frac{\Gamma(n+1/2)}{\Gamma(n+1/2)} \left[ \frac{\Gamma(n+1/2)}{\rho_{\nu}^{n+1}} - \frac{\rho_{\nu}^{2}}{\rho_{\nu}^{n+1}} - \frac{\rho_{\nu}^{2}}{\rho_{\nu}^{n+1}} \right]_{\alpha}$			
M2L	$\frac{\partial}{\partial h_{adl}} \sum_{H=0}^{H_{adl}} \frac{i^n \pi^{n-U^2} Y_n^m (\hat{V}_H) v_H^{n-2} \exp\left(-\pi^2 v_H^2 / \kappa^2\right)}{V \Gamma(n+1/2)} = -\frac{i^n \pi^{n-U^2} v_H^{n-4}}{V \Gamma(n+1/2)} \exp\left(-\frac{\pi^2 v_H^2}{\kappa^2}\right)$			
(image cells)	$\frac{\partial h_{\alpha\beta}}{\partial h_{\mu\neq0}} \frac{2}{V\Gamma(n+1/2)} = -\frac{1}{V\Gamma(n+1/2)} \exp\left(-\frac{1}{\kappa^2}\right)$			
	$\times \sum_{H=0}^{H_{mn}} \left\{ Y_{s}^{m} \left( \hat{v}_{H} \right) v_{H}^{2} h_{\beta \alpha}^{-1} + \sum_{\gamma} v_{H\alpha} h_{\beta \gamma}^{-1} \left[ Y_{s}^{m} \left( \hat{v}_{H} \right) \left\{ (n-2) - \frac{2\pi^{2} v_{H}^{2}}{\kappa^{2}} \right\} v_{H\gamma} - \left( v_{H} \times i \hat{l}_{H} \right)_{\gamma} Y_{s}^{m} \left( \hat{v}_{H} \right) \right] \right]$			
M2L	$\partial \stackrel{\mathbf{v}_{m}}{\longrightarrow} Y_{m,i}^{m-i}\left(\hat{\mathbf{p}}_{v}\right) \stackrel{\mathbf{v}_{m}}{\longrightarrow} Y_{i}^{m}\left(\hat{\mathbf{p}}_{v}\right) \Gamma\left(n+1/2, \kappa^{2} \rho_{v}^{2}\right)  \partial \stackrel{\mathcal{H}_{m}}{\longrightarrow} i^{n} \pi^{n+1/2} Y_{i}^{m}\left(\hat{v}_{H}\right) v_{H}^{n-2} \exp\left(-\pi^{2} v_{H}^{2} / \kappa^{2}\right)$			
(image cells)	$\frac{\partial}{\partial h_{\alpha\beta}} \sum_{\nu=0}^{V_{me}} \frac{Y_{n+j}^{m+\ell}(\hat{\boldsymbol{\rho}}_{\nu})}{\rho_{\nu}^{n+j+1}} = \sum_{\nu=0}^{V_{m}} \frac{Y_{n}^{m}(\hat{\boldsymbol{\rho}}_{\nu})}{\rho_{\nu}^{n+1}} \frac{\Gamma(n+1/2,\kappa^{2}\rho_{\nu}^{2})}{\Gamma(n+1/2)} + \frac{\partial}{\partial h_{\alpha\beta}} \sum_{H=0}^{U_{me}} \frac{i^{n}\pi^{n-1/2}Y_{n}^{m}(\hat{\boldsymbol{\nu}}_{H})\nu_{H}^{n-2}\exp(-\pi^{2}\nu_{H}^{2}/\kappa^{2})}{V\Gamma(n+1/2)}$			
M2L	$\frac{\partial L_{j}^{k}(0)}{\partial L_{j}} \sum_{n=1}^{\infty} \sum_{k,w}^{n} \left[ \frac{\partial M_{n}^{m}(0)}{\partial \mathbf{v}_{n}} Y_{n+j}^{w-k}(\hat{\mathbf{p}}_{\nu}) + \mathbf{M}_{n}^{m}(0) \left[ \frac{\partial}{\partial \mathbf{v}_{n}} Y_{n+j}^{w-k}(\hat{\mathbf{p}}_{\nu}) \right] \right]$			
(image cells)	$\frac{\partial L_{j}^{k}(0)}{\partial h_{\alpha\beta}} = \sum_{n=0}^{\infty} \sum_{m=-n}^{n} T_{\mathrm{ML},j,n} \left\{ \frac{\partial M_{n}^{m}(0)}{\partial h_{\alpha\beta}} \sum_{\nu=0}^{\nu_{m}} \frac{Y_{n+j}^{m-k}(\hat{p}_{\nu})}{\rho_{\nu}^{n+j+1}} + M_{n}^{m}(0) \left[ \frac{\partial}{\partial h_{\alpha\beta}} \sum_{\nu\neq0}^{\nu_{m}} \frac{Y_{n+j}^{m-k}(\hat{p}_{\nu})}{\rho_{\nu}^{n+j+1}} \right]_{\alpha} s_{\nu\beta} \right\}$			
L2L	$\frac{\partial L_{j}^{k}\left(l+1\right)}{\partial h_{\alpha\beta}} = \sum_{n=j}^{p} \sum_{m=-n}^{n} T_{\mathrm{LL},j,n} \left\{ \frac{\partial L_{n}^{m}\left(l\right)}{\partial h_{\alpha\beta}} Y_{n-j}^{m-k}\left(\hat{\boldsymbol{\rho}}_{\mathrm{LL}}\right) \boldsymbol{\rho}_{\mathrm{LL}}^{n-j} + L_{n}^{m}\left(l\right) \left[ \nabla_{\mathrm{LL}} Y_{n-j}^{m-k}\left(\hat{\boldsymbol{\rho}}_{\mathrm{LL}}\right) \boldsymbol{\rho}_{\mathrm{LL}}^{n-j} \right]_{\alpha} s_{\mathrm{LL},\beta} \right\}$			
L2P	$P_{\alpha\beta} = -\frac{1}{V} \left\{ \frac{1}{2} \sum_{j=0}^{P} \sum_{k=-j}^{j} \sum_{\gamma} \frac{\partial L_{j}^{k} \left( l_{\max} \right)}{\partial h_{\alpha\gamma}} h_{\gamma\beta} \sum_{l=1}^{N} q_{l} Y_{j}^{k} \left( \hat{r}_{l} \right) r_{l}^{j} - \frac{1}{2} \sum_{l=1}^{N} f_{l\alpha} r_{l\beta} \right\} + P_{\text{polags}}$			

Work by Prof. N. Yoshii cf. Isotropic pressure (implemented in present MODYLAS 1.0.4)

$$P_{\rm FMM} = \frac{\Phi_{\rm FMM}}{3V}$$



Parrinello-Rahman NPT ensemble MD

⇒anisotropic systems, e.g. crystal, membrane...

#### Physical property

⇒interfacial / surface tension ...

#### [7] N. Yoshii, Y. Andoh, S. Okazaki, J. Comput. Chem., in press.

# Conclusion

•Our group have developed a general-purpose MD calculation software MODYLAS combined with FMM.

•MODYLAS has been creating remarkable output of scientific researches with the K computer, and its successor models (FX100).

• To realize large-scale MD calculation on the post K computer, we continue to make many efforts to improve both performance of arithmetic operation and MPI communication, and to extend functions of software.

•We need further efforts to achieve the goal

## How to download MODYLAS



#### Full version of MODYLAS :

MODYLAS MODYLAS MOlecular DYnamics software for LAS HOME OVERVIEW DOWNLOAD DOCUMENTATION RELEASE NOT HOME • MODYLAS	TE FORUMS LITERATURE DEVELOPERS CONTACT LINKS	http://fiber-miniapp.github	
MODYLAS VERSION: 0.9.0beta COMMENT: work on the K-Compute If you agree our license, please write your email address to send a URL for download, and the	Manual(PDF) Tutorial(PDF)	<b>FIBER</b>	
Your personal information filled out boow will be used to send MODYLAS news and report a For use of MODYLAS, please cite the following: "MODYLAS: A Highly Parallelized C meral-Purpose Molecular Dynamics Simulation Progra Highly Scalable Fine-Grained New Pa allel Processing Algorithms", J. Chem. Theo. Comp., 9		Code is open, but it is not "open source". Read LICENSE.pdf in package.	
Multiple Sector Registration	You will receive e-mail #出人 modylas-admin@draco.ims.ac.jp論 作名 Download link for http://www.modylas.org/ 宛先 yoshimichi.andoh@apchem.nagoya-u.ac.jp常 Dear visitor, Thank you for your interest. Please use the following link to download http://www.modylas.org/node/19/download/ This link will be accessible until Wed, access after the link expires, don't hes http://www.modylas.org/	a233f49281a202d6de3ca1a9ccfd2538 01/22/2014 - 13:31. If you need	

# **Related publications**

#### Software developments:

[1] Y. Andoh, N. Yoshii, K. Fujimoto, et al., "MODYLAS: A highly parallelized general-purpose molecular dynamics simulation program for large-scale systems with long-range forces calculated by fast multipole method (FMM) and highly scalable fine-grained new parallel processing algorithms", *J. Chem. Theory Comput.*, **2013**, 9, 3201.

[2] Y. Andoh, S. Suzuki, S. Ohshima, T. Sakashita, M. Ogino, T. Katagiri, N. Yoshii, and S. Okazaki, "A thread-level parallelization of pairwise additive potential and force calculations suitable for current many-core architectures", *J. of Supercomputing*, in press
[3] N. Yoshii, Y. Andoh, and S. Okazaki, "Pressure tensor for electrostatic interaction calculated by fast multipole method with periodic boundary condition", *J. Comput. Chem.*, in press.
[4] Y. Andoh, N. Yoshii, and S. Okazaki, "Fast multipole method adaptive to a rectangular unit cell and partitioned tree structure with a mixture of binary and ternary branches", *will be submitted soon*.

#### **Scientific outputs:**

[5] Y. Andoh, N. Yoshii, A. Yamada, K. Fujimoto, H. Kojima, M. Mizutani, A. Nakagawa, A. Nomoto, S. Okazaki, "All-atom molecular dynamics simulation study of entire polivirus empty capsids in solution", J. Chem. Phys., 141, 165101 (2014).

[6] T.Yagasaki, M.Matsumoto, Y.Andoh, S.Okazaki, H.Tanaka, "Effect of bubble formation on the dissociation of methane hydrate in water: A molecular dynamics study", *J. Phys. Chem. B*, **118**, 1900 (2014).

[7] T.Yagasaki, M.Matsumoto, Y.Andoh, S.Okazaki, H.Tanaka, "Dissociation of methane hydrate in aqueous NaCl solutions", *J. Phys. Chem. B*, **118**, 11797 (2014).;